

Height representation of XOR-Ising loops via bipartite dimers

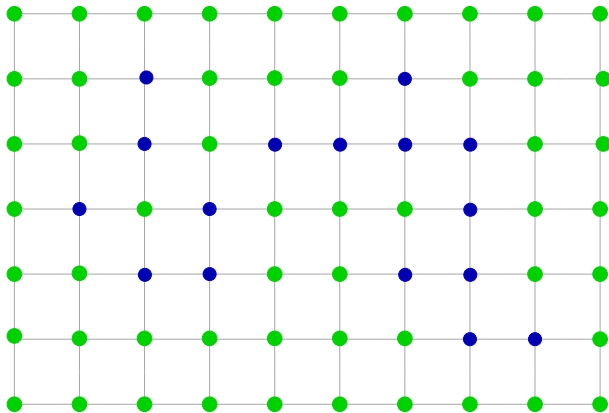
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LAGA Université Paris Nord – March 11, 2015

The Ising model and the XOR-Ising model

The Ising model

- ▶ Let $G = (V, E)$ be a finite graph embedded in the plane
- ▶ **spin configuration** $\sigma : V \rightarrow \{-1, +1\}$
- ▶ σ assigns to every vertex x a spin $\sigma_x \in \{-, +\}$



+1/-1 are represented by green/blue dots.

The Ising model

- ▶ Edges of G are assigned positive **coupling constants**:

$$J = (J_e)_{e \in E}.$$

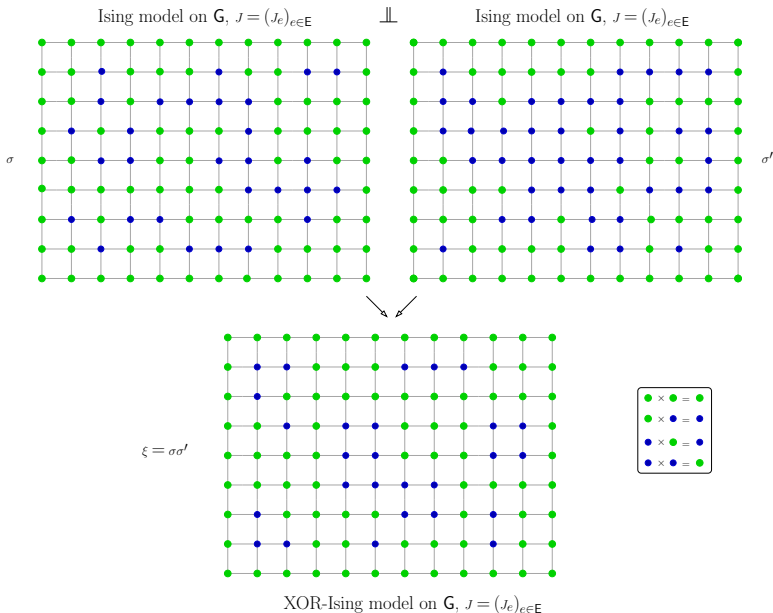
- ▶ **Ising Boltzmann measure**:

$$\forall \sigma \in \{-1, 1\}^V, \quad \mathbb{P}_{\text{Ising}}(\sigma) = \frac{1}{Z_{\text{Ising}}(G, J)} \exp \left(\sum_{e=xy \in E} J_{xy} \sigma_x \sigma_y \right),$$

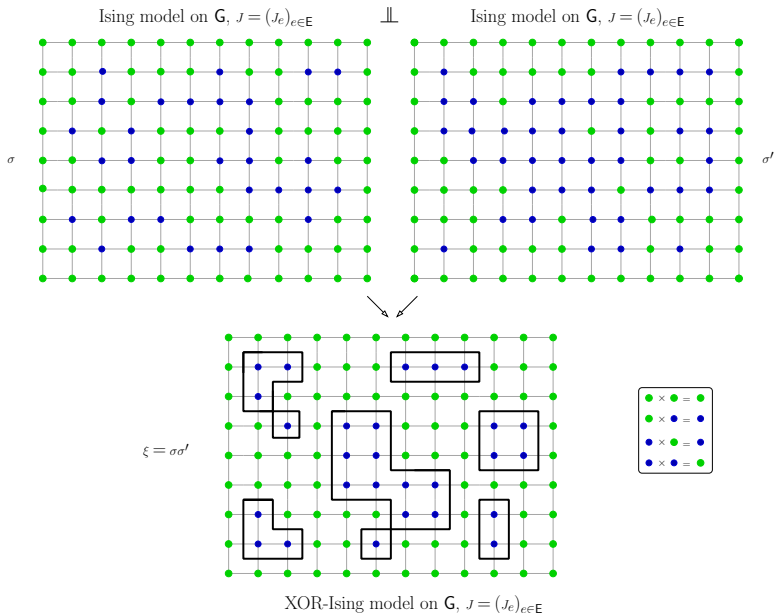
where $Z_{\text{Ising}}(G, J) = \sum_{\sigma \in \{-1, 1\}^V} \exp \left(\sum_{e=xy \in E} J_{xy} \sigma_x \sigma_y \right)$ is the

Ising partition function.

The XOR-Ising model



The XOR-Ising model



Conjecture for the XOR-Ising model

Conjecture (Wilson (11), Ikhlef–Picco–Santachiara)

The scaling limit of polygon configurations separating ± 1 clusters of the critical XOR-Ising model are contour lines of the Gaussian free field, with the heights of the contours spaced $\sqrt{2}$ times as far apart as they are for [...] the double dimer model on the square lattice.

Theorem (B–dT)

- ▶ *Polygon configurations of the XOR-Ising model have the same law as a family of contours in a bipartite dimer model.*
- ▶ *This family of contours are the level lines of a restriction of the height function of this bipartite dimer model.*

Remark

Proved when the graph G is embedded in a surface of genus g , or when G is planar, infinite.

- ▶ When the XOR-Ising model is **critical**, so is the bipartite dimer model.
- ▶ Using results of [dT] on the convergence of the height function, this gives partial proof of Wilson's conjecture.

Contour expansion of the Ising partition function [Kramers & Wannier]

Low temperature expansion

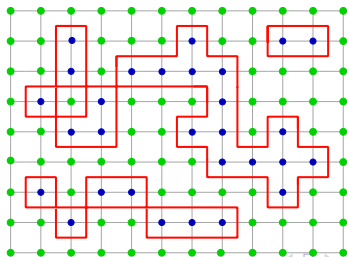
- ▶ **Polygon configuration:** subset of edges s.t. each vertex is incident to an even number of edges.

- ▶ Write
$$e^{J_e \sigma_x \sigma_y} = e^{J_e} (\delta_{\{\sigma_x = \sigma_y\}} + e^{-2J_e} \delta_{\{\sigma_x \neq \sigma_y\}}).$$

The partition function is then equal to **(LTE)**:

$$Z_{\text{Ising}}(\mathbb{G}, J) = \sum_{\sigma \in \{-1, 1\}^V} \prod_{e=xy \in E} e^{J_e \sigma_x \sigma_y} = \mathcal{C} \sum_{\mathbb{P}^* \in \mathcal{P}(\mathbb{G}^*)} \prod_{e^* \in \mathbb{P}^*} e^{-2J_e}.$$

- ▶ Geometric interp: polygon config. separate clusters of ± 1 spins.



High temperature expansion

- ▶ Write, $e^{J_e \sigma_x \sigma_y} = \cosh(J_e)(1 + \sigma_x \sigma_y \tanh(J_e))$.

The partition function is then equal to **(HTE)**:

$$Z_{\text{Ising}}(\mathbf{G}, J) = \sum_{\sigma \in \{-1, 1\}^V} \prod_{e=xy \in E} e^{J_e \sigma_x \sigma_y} = \mathcal{C}' \sum_{\mathbf{P} \in \mathcal{P}(\mathbf{G})} \prod_{e \in \mathbf{P}} \tanh(J_e).$$

- ▶ No geometric interpretation using spin configurations.

Mixed contour expansion for the double Ising model

The double Ising model

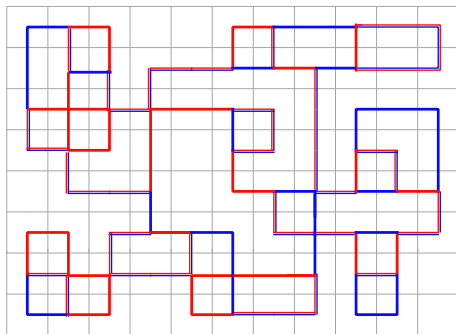
- ▶ Take 2 independent copies (red/blue) of an Ising model on G , with coupling constants J .
- ▶ Using the LTE, consider the probability measure $\mathbb{P}_{2\text{-Ising}}$: if P^* , P^* are two polygon configurations.

$$\mathbb{P}_{2\text{-Ising}}(P^*, P^*) = \frac{\mathcal{C}^2 \left(\prod_{e^* \in P^*} e^{-2J_e} \right) \left(\prod_{e^* \in P^*} e^{-2J_e} \right)}{Z_{2\text{-Ising}}(G, J)},$$

where $Z_{2\text{-Ising}}(G, J) = Z_{\text{Ising}}(G, J)^2$.

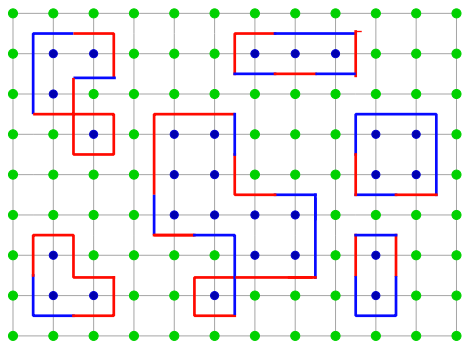
The double Ising model

- ▶ Let P^* , P^* be two polygon configurations.
- ▶ Consider the superimposition $P^* \cup P^*$.



- ▶ Define two new edge configurations:
 - ▶ $\text{Mono}(P^*, P^*)$: monochromatic edges.
 - ▶ $\text{Bi}(P^*, P^*)$: bichromatic edges.

Monochromatic edges



Monochromatic edge configuration of $P^* \cup P^*$

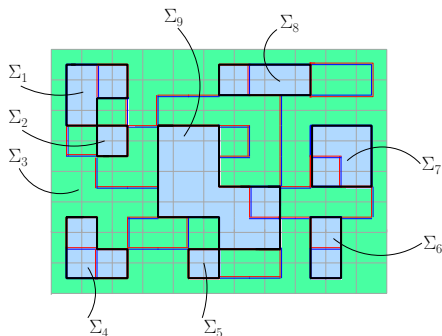
Lemma

$\text{Mono}(P^*, P^*)$ is the polygon configuration separating ± 1 clusters of the corresponding XOR-Ising spin configuration.

Goal: understand the law of monochromatic edge configurations.

Bichromatic edge configurations

- ▶ Let (P^*, P^*) be two polygon configurations.
- ▶ $\text{Mono}(P^*, P^*)$ splits the surface into connected comp. $(\Sigma_i)_i$.



Lemma

For every i , the restriction of $\text{Bi}(P^*, P^*)$ to Σ_i is the LTE of an Ising configuration on G_{Σ_i} , with coupling constants $(2J_e)$.

Probability of monochromatic configurations

Lemma

Let P^* be a polygon configuration, separating the surface into n connected components. For every i , let P_i^* be a polygon configuration of $G_{\Sigma_i}^*$.

Then, there are 2^n pairs of polygon configurations (P^*, P^*) having P^* as monochromatic edges, and P_1^*, \dots, P_n^* as bichromatic edges.

Denote by $W(P^*)$ the contribution to $Z_{2\text{-Ising}}(G, J)$ of the pairs of polygon configurations (P^*, P^*) such that $\text{Mono}(P^*, P^*) = P^*$.

Corollary

- ▶ $W(P^*) = \mathcal{C}(\prod_{e^* \in P^*} e^{-2J_e}) \prod_{i=1}^n (2Z_{\text{LT}}(G_{\Sigma_i}^*, 2J))$
- ▶ $Z_{2\text{-Ising}}(G, J) = \sum_{P^* \in \mathcal{P}(G^*)} W(P^*)$

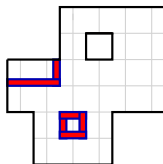
$$\mathbb{P}_{2\text{-Ising}}(\text{Mono} = P^*) = \frac{W(P^*)}{Z_{2\text{-Ising}}(G, J)}$$

Mixed contour expansion

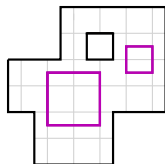
$$W(P^*) = \mathcal{C}(\prod_{e^* \in P^*} e^{-2J_e}) \prod_{i=1}^n \left(2Z_{\text{LT}}(\mathbf{G}_{\Sigma_i}^*, 2J) \right).$$

Idea [Nienhuis]: high temperature expansion in each connected component Σ_i .

$$Z_{\text{LT}}(\mathbf{G}_{\Sigma_i}^*, 2J) = \mathcal{C}(\Sigma_i) Z_{\text{HT}}(\mathbf{G}_{\Sigma_i}, 2J).$$



Low temp. expansion on $\mathbf{G}_{\Sigma_i}^*$



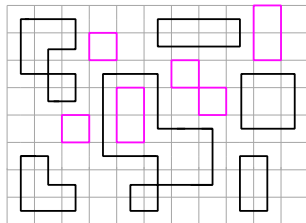
High temp. expansion on \mathbf{G}_{Σ_i} .

Mixed contour expansion

Proposition

For every polygon configuration P^* ,

$$W(P^*) = \mathcal{C} \prod_{e^* \in P^*} \left(\frac{2e^{-2J_e}}{1 + e^{-4J_e}} \right) \sum_{\{P \in \mathcal{P}(G) : P^* \cap P = \emptyset\}} \prod_{e \in P} \left(\frac{1 - e^{-4J_e}}{1 + e^{-4J_e}} \right)$$



$$\mathbb{P}_{2\text{-Ising}}(\text{Mono} = P^*) = \frac{\prod_{e^* \in P^*} \left(\frac{2e^{-2J_e}}{1 + e^{-4J_e}} \right) \sum_{\{P \in \mathcal{P}(G) : P^* \cap P = \emptyset\}} \prod_{e \in P} \left(\frac{1 - e^{-4J_e}}{1 + e^{-4J_e}} \right)}{\sum_{P^* \in \mathcal{P}(G^*)} \dots}$$

Higher genus

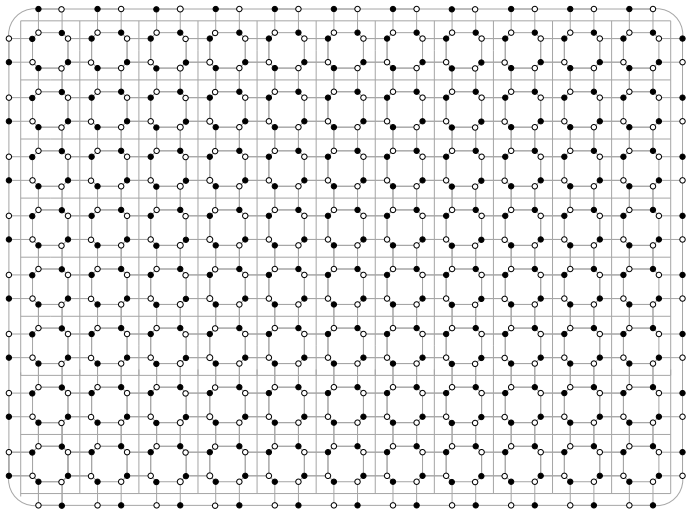
If the graph is embedded in a surface Σ of genus $g \geq 0$.

- ▶ Consider $H_1(\Sigma, \mathbb{Z}/2\mathbb{Z}) \simeq \{0, 1\}^{2g}$.
- ▶ Family of Ising models, indexed by $\varepsilon \in \{0, 1\}^{2g}$.
- ▶ The double Ising model partition function is defined as:

$$Z_{2\text{-Ising}}(\mathbf{G}, J) = \sum_{\varepsilon \in \{0, 1\}^{2g}} Z_{\text{Ising}}^{\varepsilon}(\mathbf{G}, J)^2.$$

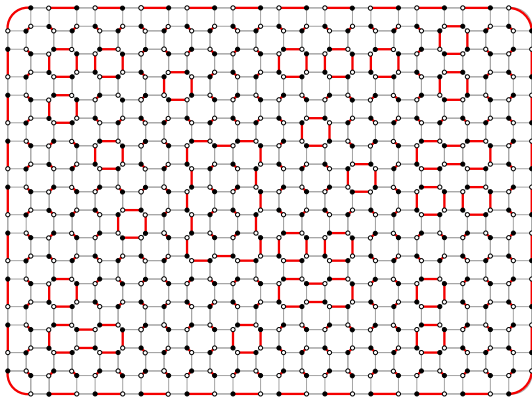
From mixed polygon configurations to dimers

The graph $G^Q = (V^Q, E^Q)$



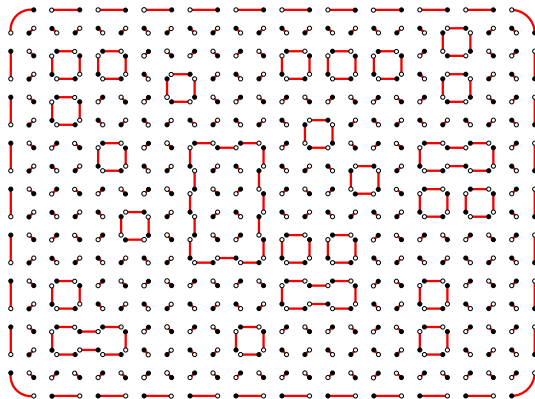
The dimer model on G^Q

dimer configuration of G^Q : a subset of edges M such that each vertex is incident to exactly one edge of M



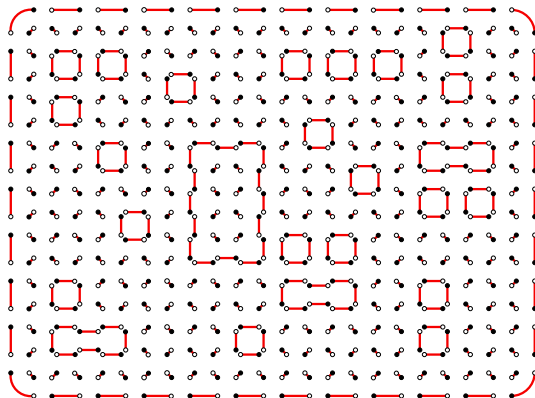
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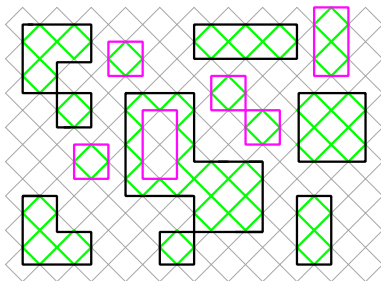
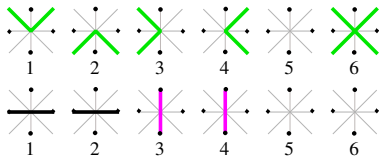


weight function ν on the edges

Dimer Boltzmann measure: $\mathbb{P}_{\text{dimer}}(M) \propto \prod_{e \in EQ} \nu_e$

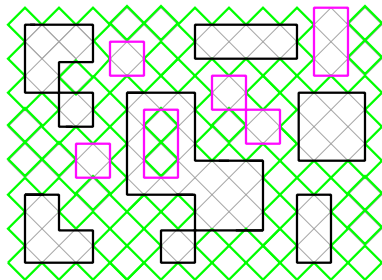
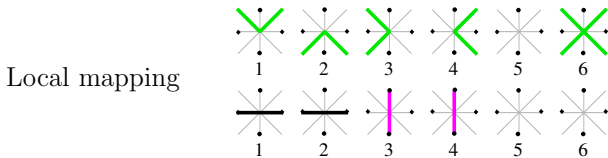
First step: from polygons to 6-vertex [Nienhuis]

Local mapping



$$\text{Weights: } \omega_{12} = \frac{2e^{-2J_e}}{1+e^{-4J_e}}, \omega_{34} = \frac{1-e^{-4J_e}}{1+e^{-4J_e}}, \omega_{56} = 1.$$

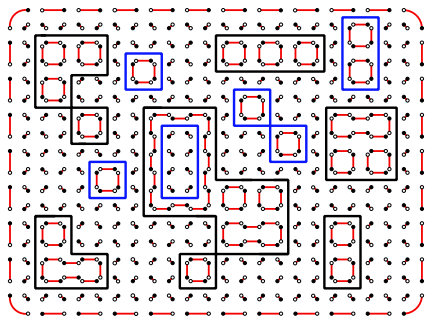
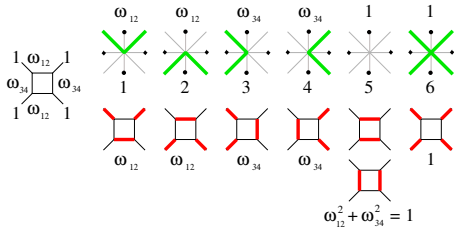
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Second step: from 6V to dimers [Wu-Lin, Dubédat]

Local mapping



Conclusion

- ▶ To every dimer configuration M of G^Q , assign

$$\text{Poly}(M) = (\text{Poly}_1(M), \text{Poly}_2(M)),$$

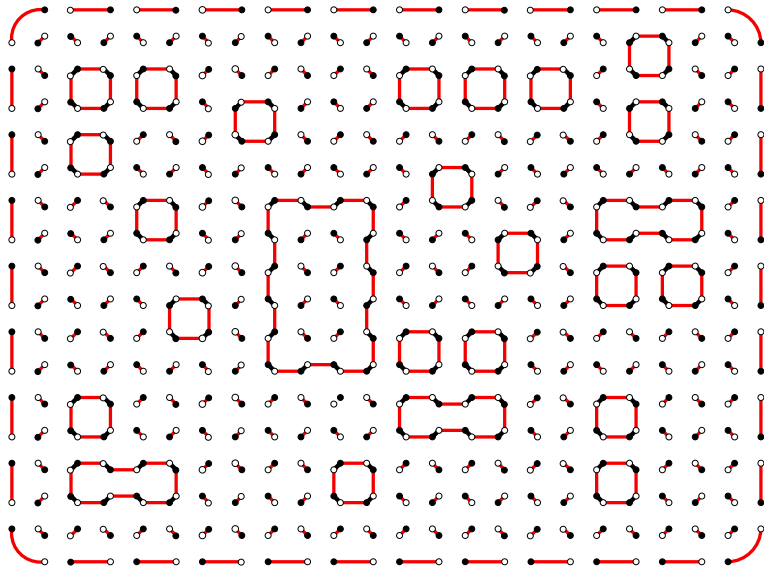
the pair of polygon configurations given by the mappings.

Theorem

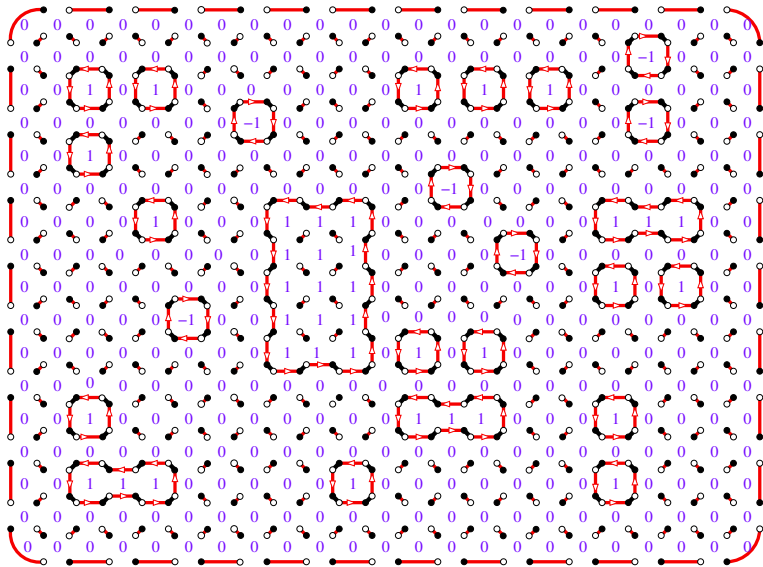
For every polygon configuration P^ of G^* ,*

$$\mathbb{P}_{2\text{-Ising}}(\text{Mono} = P^*) = \mathbb{P}_{\text{dimer}}(\text{Poly}_1 = P^*)$$

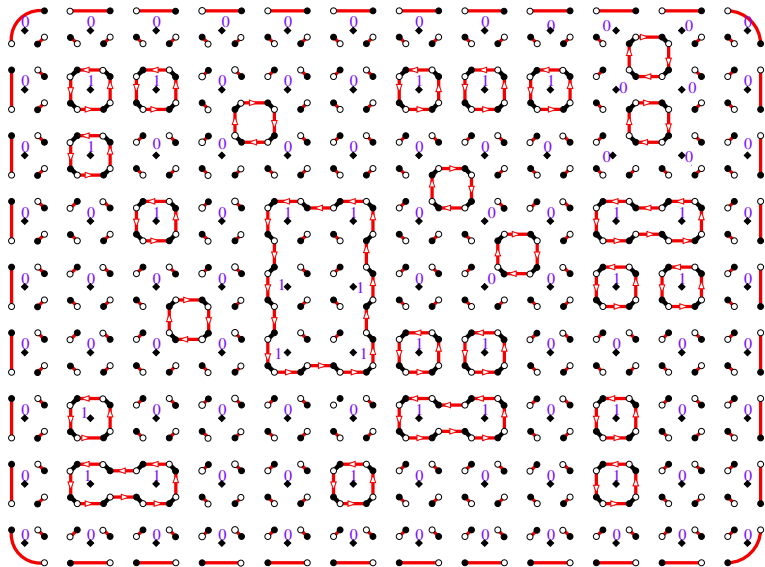
Height function for bipartite dimers (Thurston)



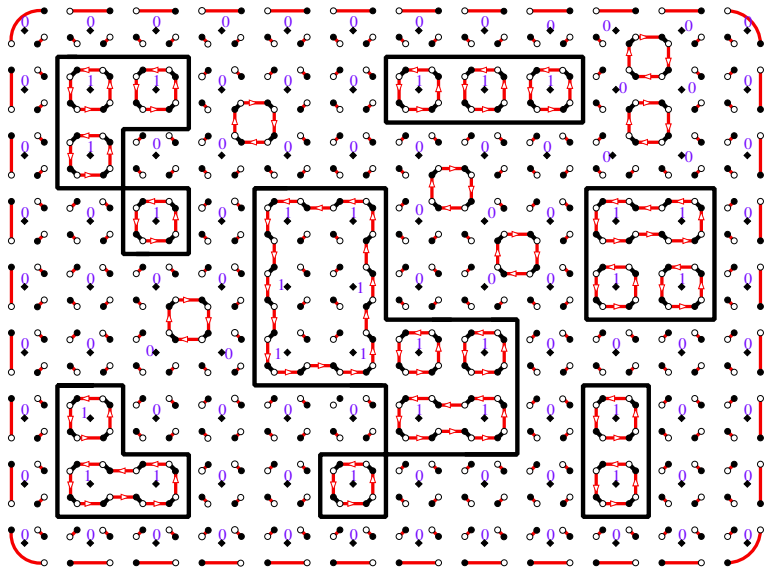
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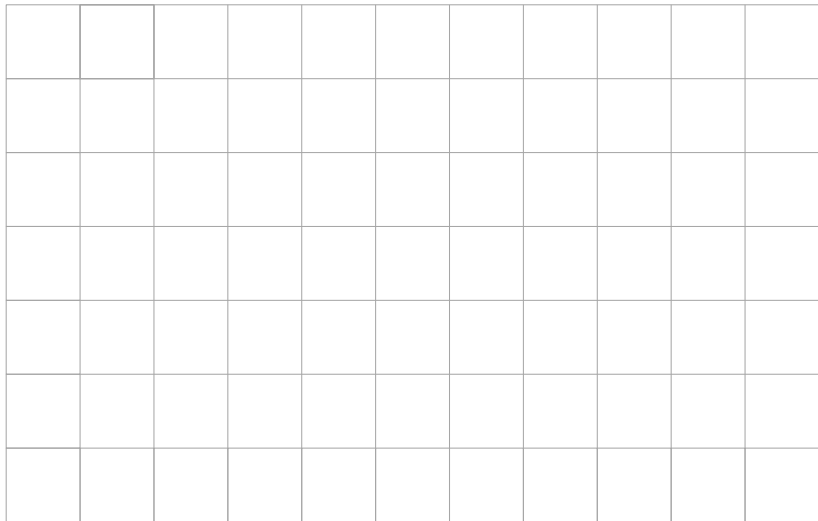


The critical XOR-Ising model on isoradial graphs

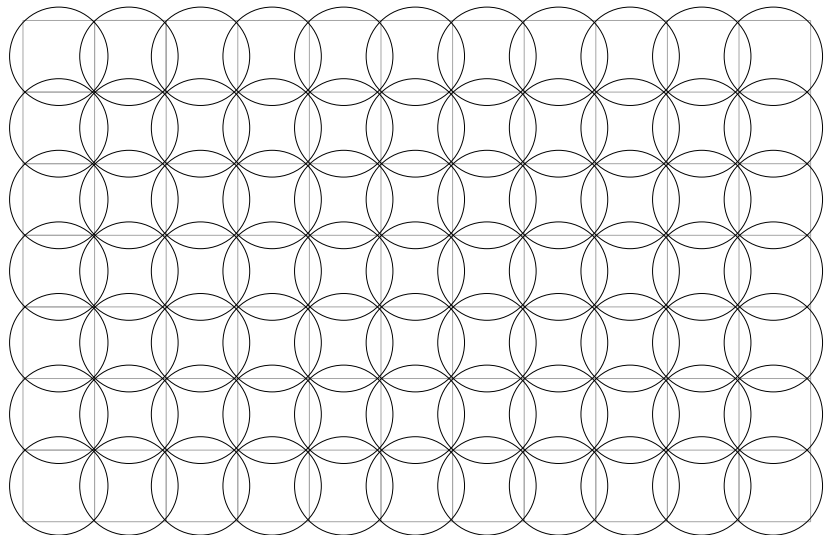
Isoradial graphs

A graph G is **isoradial** if it is planar and can be embedded in the plane in such a way that all faces are inscribed in a circle of radius 1, and that the circumcenters are in the interior of the faces.
[Duffin, Mercat, Kenyon]

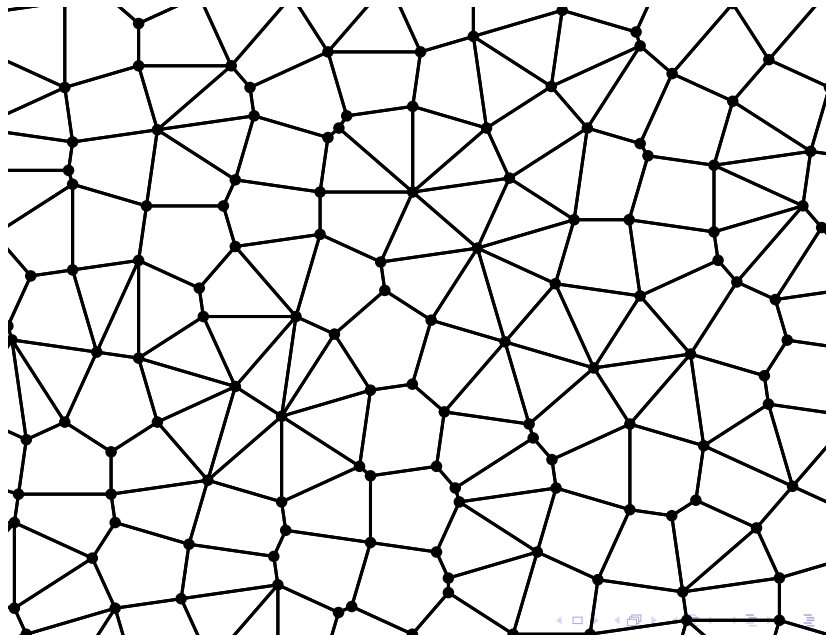
Isoradial graphs



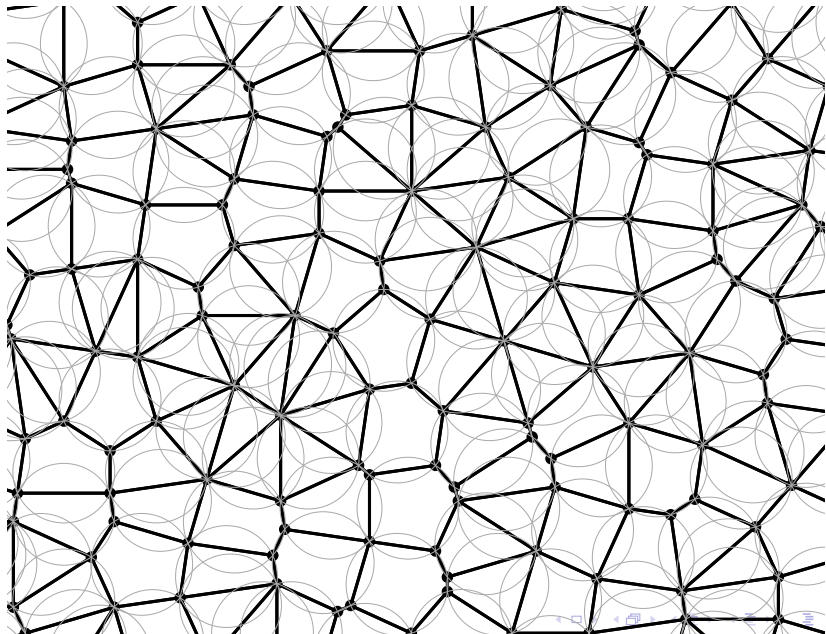
Isoradial graphs



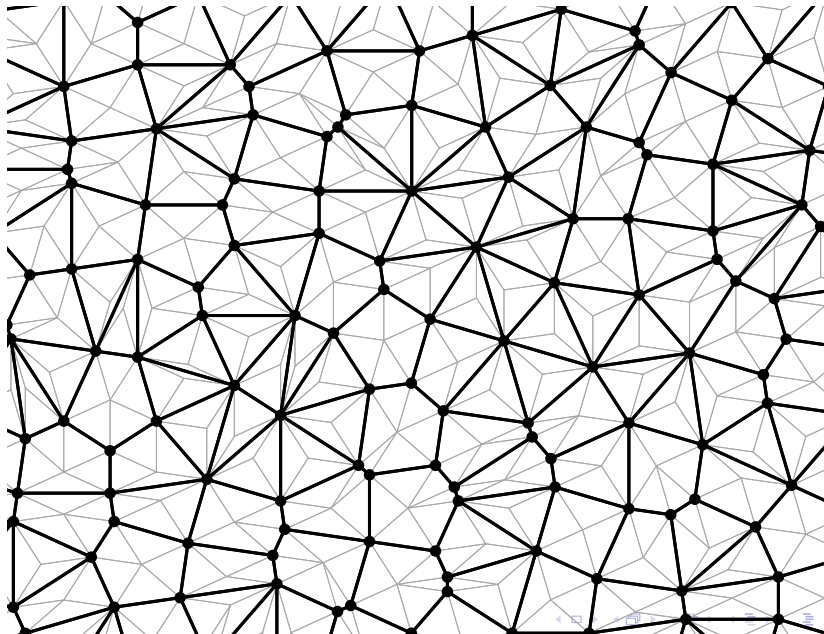
Isoradial graphs



Isoradial graphs



Associated rhombus graph



Critical Ising model on isoradial graphs

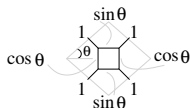
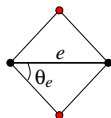
- ▶ To each edge e is naturally associated an angle θ_e
- ▶ The Ising model defined on an isoradial graph G is **critical** if the coupling constants are given by:

$$J_e = \frac{1}{2} \log \left(\frac{1 + \sin \theta_e}{\cos \theta_e} \right).$$

(\mathbb{Z} -invariance + duality [Baxter], proof in period. case [Li, Duminil-Cimasoni])

Example: $G = \mathbb{Z}^2$: $\theta_e = \frac{\pi}{4}$, $J_e = \frac{1}{2} \log(1 + \sqrt{2})$.

- ▶ The corresponding bipartite graph G^Q is also isoradial, and the weights are the **critical** dimer weights:



Back to Wilson's conjecture

Conjecture (Wilson)

The scaling limit of polygon configurations separating ± 1 clusters of the critical XOR-Ising model are contour lines of the Gaussian free field, with the heights of the contours spaced $\sqrt{2}$ times as far apart as they are for [...] the double dimer model on the square lattice.

Theorem (B-dT)

XOR-polygon configurations of the double Ising model on G have the same law as level lines of a restriction of the height function of the bipartite dimer model on $G^{\mathbb{Q}}$, with an explicit coupling.

Theorem (dT)

The height function (as a random distribution) of the critical dimer model defined on a bipartite graph converges weakly in law to $\frac{1}{\sqrt{\pi}}$ a Gaussian free field of the plane.

Back to Wilson's conjecture

Suppose we had strong form of convergence, allowing for convergence of level lines. Then:

level lines of h^ε	\rightarrow	level lines of GFF	
$(k, k \in \mathbb{Z})$		$(\sqrt{\pi}k, k \in \mathbb{Z})$	
$(k + \frac{1}{2}, k \in \mathbb{Z})$		$(\frac{\sqrt{\pi}}{2}(2k + 1), k \in \mathbb{Z})$	XOR loops

For the critical double dimer model. The height function is $h_1^\varepsilon - h_2^\varepsilon$, where h_1 and h_2 are independent, and each converges weakly in distribution to $\frac{1}{\sqrt{\pi}}$ a Gaussian free field. Thus, $h_1 - h_2$ converges weakly in distribution to $\frac{\sqrt{2}}{\sqrt{\pi}}$ a Gaussian free field.

level lines of $h_1^\varepsilon - h_2^\varepsilon$	\rightarrow	level lines of GFF	
$(k, k \in \mathbb{Z})$		$(\frac{\sqrt{\pi}}{\sqrt{2}}k, k \in \mathbb{Z})$	
$(k + \frac{1}{2}, k \in \mathbb{Z})$		$(\frac{\sqrt{\pi}}{2\sqrt{2}}(2k + 1), k \in \mathbb{Z})$	d-dimer loops