Differentially algebraic equations in physics Youssef Abdelaziz, Jean-Marie Maillard (Université Paris VI) Based on "Modular forms, Schwarzian conditions, and symmetries of differential equations in physics", arXiv 1611.08493

Séminaire CALIN

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#### Hamiltonian of the Ising model

$$H = \sum_{j,k} \{ J_{\nu} \sigma_{j,k} \sigma_{j+1,k} + J_{h} \sigma_{j,k} \sigma_{j,k+1} \}$$



- $J_{v}$ ,  $J_{h}$ : vertical and horizontal coupling constants
- The spins take the values  $\sigma_{j,k} = \pm 1$ .
- The partition function:  $\exp(-\frac{1}{k_bT}H)$

#### Nature of power series

- Algebraic:  $S(x) \in \mathbb{Q}(x)$  root of a polynomial P(t, S(t)) = 0
- D-finite:  $S(x) \in \mathbb{Q}(x)$  satisfying a linear differential equation with polynomial coefficients  $c_r(t)S^{(t)}(t) + s + c_0(t)S(t) = 0$
- Hypergeometric:  $S(x) = \sum_{n=0}^{\infty} s_n x^n$  s.t.  $\frac{s_{n+1}}{s_n} \in \mathbb{Q}(n)$ . E.g., the Gauss hypergeometric function:

$$_{2}F_{1}([a,b],[c],x) = \sum_{n=0}^{\infty} \frac{(a)_{n}(b)_{n}}{(c)_{n}} \frac{t^{n}}{n!},$$

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$$_{2}F_{1}([a,b],[c],x) = \sum_{n=0}^{\infty} \frac{(a)_{n}(b)_{n}}{(c)_{n}} \frac{t^{n}}{n!}, \quad (a)_{n} := a(a+1)\cdots(a+n-1)$$

- E.g.:  $_{2}F_{1}(1,1;1;z) = \frac{1}{1-z}, \quad _{2}F_{1}(1,1;2;z) = -\frac{\ln(1-z)}{z}$
- Partition function 2D square Ising model [Viswanathan, 2014]

$$_{4}F_{3}([1,1,\frac{3}{2},\frac{3}{2}],[2,2,2],16k^{2}]), \quad k=rac{ anh(2eta J)}{2\cosh(2eta J)}$$

#### Magnetic susceptibility of 2D Ising model

$$\chi := \beta \sum_{n=0}^{\infty} \chi^{(2n+1)}$$



ability of a material to align itself with an external imposed magnetic field

 $\chi^{(2n+1)} \longrightarrow 2n$  multiple integrals , e.g.  $\chi^{(3)}$  is given by the double integral:

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and

#### Feynman diagrams are D-finite



Feynman diagrams  $\rightarrow$  first order perturbations of *n*-fold integral of the operator *S* (scattering operator) giving the probability of such interactions:

$$S = \sum_{n=0}^{\infty} \frac{\iota^n}{n!} \underbrace{\int \cdots \int}_{j=1}^n \frac{d^4 x_j}{d^4 x_j} T \prod_{j=1}^n L(x_j)$$

 $L_v(x_j) \longrightarrow$  Lagrangian of interaction, T the time ordered product of operators,  $d^4x_j$  four-vectors

#### Multiple integrals of an algebraic object



(D-finite = solution of linear ODE with polynomial coefficients)

#### Diagonal of a rational function

For a formal power series F given by

$$F(z_1, z_2, \cdots, z_n) = \sum_{m_1=0}^{\infty} \cdots \sum_{m_n=0}^{\infty} F_{m_1, \cdots, m_n} z_1^{m_1} \cdots z_n^{m_n},$$

the diagonal of F is defined as the single variable series:

$$Diag(F(z_1, z_2, \cdots, z_n)) := \sum_{m=0}^{\infty} F_{m, \cdots, m} z^m$$

Example. One of the many diagonals leading to Apéry numbers:

$$\mathsf{Diag}\frac{1}{(1-z_1-z_2)(1-z_3-z_4)-z_1z_2z_3z_4} = \sum_{n\geq 0}\sum_{k=0}^n \binom{n}{k}^2 \binom{n+k}{k}^2 z^n$$

### $_{2}F_{1}$ , modular forms, and physics

The Gauss hypergeometric function  $_2F_1 \rightarrow \text{PHYSICS}!$ , e.g. the differential operator of  $\chi^{2n+1}$  factorizes into operators that annihilate  ${}_2F_1$  functions. A. Bostan, S. Boukraa, S. Hassani, M. van Hoeij, J.-M. Maillard, J.-A. Weil, N. Zenine, The Ising model: from elliptic curves to modular forms and Calabi-Yau equations, 2011] [M. Assis, S. Boukraa, S. Hassani, M. van Hoeij, J.-M. Maillard, B. M. McCoy, Diagonal Ising susceptibility: elliptic integrals, modular forms and Calabi-Yau

equations, 2012

$$E_4(q) = 1 + 240 \sum_{n=0}^{\infty} n^3 \frac{q^n}{1-q^n}$$
  
=  $_2 F_1\left(\left[\frac{1}{12}, \frac{5}{12}\right], [1], \frac{1728}{j(\tau)}\right)^4$   
 $q = \exp(2i\pi\tau), j(\tau) \rightarrow j$ -invariant

#### Modular forms as pullbacked $_2F_1$ functions

- Emergence of modular forms in physics through  $_2F_1$  functions
- Modular forms emerge through covariance properties of <sub>2</sub>*F*<sub>1</sub>:

$${}_{2}F_{1}([\alpha,\beta],[\gamma],p_{1}(x)) = \mathcal{A}(x){}_{2}F_{1}([\alpha,\beta],[\gamma],p_{2}(x))$$

 $\mathcal{A}(x)$ ,  $p_1(x)$  and  $p_2(x)$  are rational functions.  $p_1(x)$  and  $p_2(x)$  are called pullbacks, the  $_2F_1$  is thus called pullbacked. For instance:

$${}_{2}F_{1}\left(\left[\frac{1}{12},\frac{5}{12}\right],\left[1\right],\frac{1728x}{(5+10x+x^{2})^{3}}\right) = \\ \left(\frac{5+10x+x^{2}}{3125+250x+x^{2}}\right)_{2}^{1/4}F_{1}\left(\left[\frac{1}{12},\frac{5}{12}\right],\left[1\right],\frac{1728x^{5}}{(3125+250x+x^{2})^{3}}\right)$$

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• w.l.o.g we have:  $\mathcal{A}(x)_2 F_1([\alpha,\beta],[\gamma],y(x)) = {}_2F_1([\alpha,\beta],[\gamma],x)$  $\mathcal{A}(x)$  and y(x) algebraic functions.

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• w.l.o.g we have:  $\mathcal{A}(x)_2 F_1([\alpha, \beta], [\gamma], y(x)) = {}_2F_1([\alpha, \beta], [\gamma], x)$  $\mathcal{A}(x)$  and y(x) algebraic functions. Modular equation M(x, y(x)) = 0:  $1953125x^3y^3 - 187500x^2y^2(x+y) + 375xy(16x^2 - 4027xy + 16y^2)$  $-64(x+y)(x^2 + 1487xy + y^2) + 110592xy = 0$ 

#### Schwarzian condition

Theorem ( Abdelaziz–Maillard, 2016 )

If we have a pullback given by:

$$\mathcal{A}(x)_{2}F_{1}([\alpha,\beta],[\gamma],x) =_{2} F_{1}([\alpha,\beta],[\gamma],y(x))$$

then we have the following "Schwarzian condition":

$$W(x) - W(y(x))y'(x)^2 + \{y(x), x\} = 0$$

where 
$$W(x) := p'(x) + \frac{p(x)^2}{2} - 2q(x)$$
  
with  $p(x) = \frac{(\alpha + \beta + 1)x - \gamma}{x(x-1)}$   $q(x) = \frac{\alpha\beta}{x(x-1)}$ 

#### NB: The Schwarzian derivative is defined by

$$\{y(x), x\} := \frac{y'''(x)}{y'(x)} - \frac{3}{2} \left(\frac{y''(x)}{y'(x)}\right)^2$$



Differentially algebraic equations in physics

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### NB: The Legendre derivative is defined by $\{y(x), x\} := \frac{y'''(x)}{v'(x)} - \frac{3}{2} \left(\frac{y''(x)}{v'(x)}\right)^2$



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Differentially algebraic equations in physics

#### Proof of our theorem of the Schwarzian condition

We introduce

• the operator  $L_2 := D_x^2 + p(x)D_x + q(x)$  annihilating  $F(x) := {}_2F_1([\alpha, \beta], [\gamma], x)$ 

• the operator  $L_2^{(c)} := \frac{1}{v(x)} L_2 v(x)$ , i.e.

$$L_2^{(c)} = D_x^2 + \left(p(x) + 2\frac{v'(x)}{v(x)}\right)D_x + q(x) + p(x)\frac{v'(x)}{v(x)} + \frac{v''(x)}{v(x)}$$

NB:  $L_2^{(c)}$  annihilates  $\mathcal{A}(x)F(x)$  (with  $\mathcal{A}(x) = 1/v(x)$ ):  $L_2^{(c)} \frac{1}{v}F(x) = \frac{1}{v}L_2 \mathcal{V}_{k}F(x) = 0$ 

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• So, the operator annihilating F(y(x)) is

$$L_2^{(p)} = D_x^2 + \left( p(y(x))y'(x) - \frac{y''(x)}{y'(x)} \right) D_x + q(y(x))y'(x)^2$$

• When does the equality  $L_2^{(c)} = L_2^{(p)}$  hold?

#### Proof of our theorem of the Schwarzian condition

Well, identifying  $L_2^{(c)} = L_2^{(p)}$  gives us two conditions:

Condition 1: 
$$p(x) + 2\frac{v'(x)}{v(x)} = p(y(x))y'(x) - \frac{y''(x)}{y'(x)}$$

Condition 2: 
$$q(x) + p(x)\frac{v'(x)}{v(x)} + \frac{v''(x)}{v(x)} = q(y(x))y''(x)^2$$

Introducing  $w(x) := \exp\left(-\int p(x)dx\right)$ , i.e.  $p(x) = -\frac{w'(x)}{w(x)}$ , Condition 1 rewrites

$$-\frac{w'(x)}{w(x)} + 2\frac{v'(x)}{v(x)} = -\frac{y''(x)}{y'(x)} + \frac{w'(y(x))}{w(y(x))}$$

Integrating the log-derivative terms we get:

$$-\ln w(x) + 2\ln v(x) = -\ln y'(x) - \ln w(y(x))$$

Taking exponential gives

$$v(x) = \sqrt{\frac{w(x)}{w(y(x))y'(x)}}$$

inserting it in Condition 2 gives the Schwarzian condition in the theorem.

Assuming that the operator is globally nilpotent is equivalent to:

$$p(x) = -\frac{w'(x)}{w(x)}$$

The following statements are a consequence of global nilpotence:

- The Wronskian is the *n*-th root of a rational function
- The solutions of the differential equation have rational coefficients
- The *p*-curvature is a nilpotent matrix mod prime
- Global nilpotence  $\rightarrow$  rational coefficients of solutions  $\rightarrow p(x) = \frac{d}{dx} \ln w(x)$

#### Modular equation M(x,y(x))=0 and modular invariant

The *j*-invariant of the elliptic curve:

$$j(k) = 256 \frac{(1-k^2+k^4)^3}{k^4(1-k^2)^2}$$

The Landen transformation:

$$k_L = \frac{2\sqrt{k}}{1+k}$$

The transform of the elliptic invariant through  $k_L$ :

$$j(k_L) = 16 \frac{(1+14k^2+k^4)^3}{k^2(1-k^2)^4}$$

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The two corresponding Hauptmoduls (similar to a group generator):

$$x = \frac{1728}{j(k)}$$
  $y = \frac{1728}{j(k_L)}$ 

are related through the modular equation  $\tau \rightarrow 2\tau$ :

$$M(x, y) = 1953125x^{3}y^{3} - 187500x^{2}y^{2}(x + y) + 375xy(16x^{2} - 4027xy + 16y^{2})$$
$$-64(x + y)(x^{2} + 1487xy + y^{2}) + 110592xy = 0$$

For one  ${}_{2}F_{1}([a, b], [1], x)$  with two different pullbacks  $\alpha x + \cdots$   $\alpha x^{2} + \cdots$  $\alpha x^{3} + \cdots$ 

we obtain the isogenies series-solution "structure"

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 $\alpha x + \cdots$   
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we obtain the isogenies series-solution "structure"

This set of solutions is either:

- Algebraic: e.g.  $_2F_1([\frac{1}{12}, \frac{5}{12}], [1], x)$ , we recover "some" commutation like in the case of isogenies (as we will see below)
- Transcendent

# Schwarzian condition and modular forms: $\tau \rightarrow 2\tau$ and beyond

The modular form:

$$\mathcal{A}(x)_{2}F_{1}\left(\left[\frac{1}{12},\frac{5}{12}\right],\left[1\right],x\right) =_{2}F_{1}\left(\left[\frac{1}{12},\frac{5}{12}\right],\left[1\right],\frac{y(x)}{y(x)}\right)$$
(1)

- $\mathcal{A}(x)$  is an algebraic function
- y(x) is an algebraic function corresponding to the modular equation corresponding to au o 2 au

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$$y(x) = \frac{1}{1728}x^2 + \frac{31}{62208}x^3 + \frac{1337}{3359232}x^4 + \frac{349115}{1088391168}x^5 + \cdots$$

The Schwarzian condition is verified in this case with:

$$W(x) = -\frac{32x^2 - 41x + 36}{72x^2(x-1)^2}, p(x) = \frac{3x-2}{2x(x-1)}, q(x) = \frac{5}{144x(x-1)}$$

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It turns out that one can write, for the modular equations corresponding to  $\tau \rightarrow N\tau$ , the function in the form of (1) above. Thus the equation (1) above encapsulates all the modular equations corresponding to  $\tau \rightarrow N\tau$ .

The modular equation of order three  $\tau \rightarrow 3\tau$ :

 $\begin{aligned} 26214400000000x^3y^3(x+y) + 409600000x^2y^2(27x^2-45946xy+27y^2) \\ + 15552000xy(x+y)(x^2+241433xy+y^2) \\ + 729x^4 - 779997924x^3y + 1886592284694x^2y^2 - 779997924xy^3 + 729y^4 \\ + 2811677184xy(x+y) - 2176782336xy = 0 \end{aligned}$ 

has the series expansion starting in  $x^3$  and given by:

$$y(x) = \frac{x^3}{2985984} + \frac{31x^4}{71663616} + \frac{36221x^5}{82556485632} + \frac{29537101x^6}{71328803586048} + \dots$$

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has the series expansion starting in  $x^3$  and given by:

$$y(x) = \frac{x}{5159780352} + \frac{31x}{92876046336} + \frac{43909x}{106993205379072} + \cdots$$

- Except for this last series solution, the solution series corresponding to the isogenies  $\tau \rightarrow N\tau$  have the form  $ax^N + \cdots$
- The series solution corresponding to au o 3 au and au o 4 au are solution of the Schwarzian condition

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Generalizing the solution series corresponding to  $\tau \rightarrow 2\tau$  we seek solution series of the Schwarzian condition of the form  $ax^2 + \cdots$ :

$$y_2 = ax^2 + \frac{31ax^3}{36} - \frac{a(5952a - 9511)}{13824}x^4 + \cdots$$

reducing to the solution of au 
ightarrow 2 au when a=1/1728

A one-parameter family of solution-series  $bx^3 + \cdots$  for the modular equation corresponding to  $\tau \rightarrow 3\tau$ :

$$y_3 = bx^3 + \frac{31b}{24}x^4 + \frac{36221b}{27648}x^5 + \cdots$$

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Finally the one-parameter series

$$y_4 = cx^4 + \frac{31c}{18}x^5 + \frac{43909c}{20736}x^6 + \cdots$$

reduces to a previous series of the form  $x^4 + \cdots$  for  $c = 1/5159780352 = 1/1728^3$ 

- These series do not commute:  $y_i(y_j(x)) \neq y_j(y_i(x))$ .
- Composing the solution series  $y_3$  and  $y_2$  with  $d = ab^2$ :

$$y_2(y_3(x)) = dx^6 + \frac{31dx^7}{12} + \frac{59285d}{13824}x^8 + \cdots$$
  
•  $y_2(y_3(x)) = y_3(y_2(x)) \leftrightarrow ?$ 

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•  $y_2(y_3(x)) = y_3(y_2(x)) \leftrightarrow ab^2 = ba^3$ 

#### Conclusion

- The Schwarzian condition encapsulates the infinite number of modular equations au o N au.
- Strong incentive to develop more differentially algebraic tools from an algorithmic perspective : to test the non-D-finiteness of the Ising susceptibility for example!
- Strong incentive to examine further the occurence of non-linear symmetries (like the Landen transformation) in physics.



- Built to generalize the differential Galois group to non-linear ODE's and non linear functional equations having the form f(x + 1) = y(f(x)).
- Having a finite non-linear differential Galois group guarantees "some integrability" and this is guaranteed by Casale's condition:

$$\nu(y)y''(x)^2 - \nu(x) + \frac{y'''(x)}{y'(x)} - \frac{3}{2} \left(\frac{y''(x)}{y'(x)}\right)^2 = 0$$

#### Modular equations: definition through $\theta$ functions

With  $q = \exp(i\pi\tau)$ ,  $\tau = iK'/K$  the  $\theta_3$  and  $\theta_4$  functions are defined as follows:

$$heta_2 = 2q^{1/4} \prod_{n \ge 1} \left( \frac{1 - q^{4n}}{1 - q^{4n-2}} \right), \quad heta_3 = \sum_{-\infty}^{\infty} q^{n^2}, \quad heta_4 = \sum_{-\infty}^{\infty} (-1)^n q^{n^2}$$

where  $K = (\pi/2)\theta_3^2(\tau)$  and  $K'(\tau) = K(-1\tau)$ . We can write the identity:

$$heta_3( au)^2 + heta_4( au)^2 = 2 heta_3(2 au)^2 = rac{2}{1+k'}$$

with  $\sqrt{k(\tau)} = \frac{\theta_2(\tau)}{\theta_3(\tau)}$ ,  $\sqrt{k'(\tau)} = \frac{\theta_4(\tau)}{\theta_3(\tau)}$  and  $l'(\tau) = k'(p\tau)$  where p is given by a positive integer, we have:

$$\frac{1}{l'} = \frac{1}{2}(\sqrt{k'} + \frac{1}{\sqrt{k'}})$$

giving in the case p = 2 the modular equation that sends  $\tau$  to  $2\tau$ .

#### Painlevé equations

- The hypergeometric function, the Bessel function, the Airy function, the Hermite polynomials, are all "special" (appearing in problems elated to physics) functions solution of linear differential equations.
- Elliptic functions are also "special" functions: they appear in physics as we shall see here, yet they are solution of simple, yet non-linear differential equations.
- Painlevé was set out to find special functions satisfying non-linear differential equations, yet have nice properties (all their singularities are poles).

Painlevé wanted to classify all differential equations of order two having the form:

$$u_{xx} = R(x, u, u_x)$$

with R being a rational function. Painlevé found 50 equations having this form, six of these were irreducible to known functions; they are known today as the six Painlevé equations.

### Magnetic susceptibility = ratio of D-finite functions?

The hypergeometric function:

$$_{2}F_{1}([1/3, 1/3], [1], 27x)$$

is D-finite and verifies the following linear differential equation

$$(27x^2-x)\left(\frac{d^2}{dx^2}F(x)\right)+(45x-1)\left(\frac{d}{dx}F(x)\right)+3F(x).$$

Similarly the hypergeometric function given by

 $_{2}F_{1}([1/2,1/2],[1],16x)$ 

verifies the **D-finite** equation

$$(16x^2-x)\left(\frac{d^2}{dx^2}F(x)\right)+(32x-1)\left(\frac{d}{dx}F(x)+4F(x)\right)$$

**Reminder**: A function is D-finite when it is solution of a *linear* differential equation and with *rational* coefficients in *x*.

#### Magnetic susceptibility = ratio of D-finite functions?

The ratio of these two D-finite functions is given by:

 $\frac{{}_2F_1([1/3,1/3],[1],27x)}{{}_2F_1([1/2,1/2],[1],16x)}$ 

- While the product of two D-finite functions is always D-finite, the ratio of two D-finite functions is generally **not so** (except if the D-finite function at the denominator is an algebraic function)!
- In fact the differential equation that this ratio verifies is non-linear as we can see in the next slide

$$\begin{aligned} -2x^{2}(27x-1)(-1+16x)((27x-1)(-1+16x)\frac{d}{dx}F(x)) \\ &-72xF(x)-F(x))\frac{d^{3}}{dx^{3}}F(x) \\ &+3x^{2}(27x-1)^{2}(-1+16x)^{2}\left(\frac{d^{2}}{dx^{2}}F(x)\right)^{2} \\ &-2x(93312\frac{d}{dx}F(x)x^{4}-7992\frac{d}{dx}F(x)x^{3} \\ &-93312x^{3}F(x)+87\frac{d}{dx}F(x)x^{2}+168x^{2}F(x) \\ &+3\frac{d}{dx}F(x)x+297xF(x)-4F(x))\frac{d^{2}}{dx^{2}}F(x) \\ &+(-1+16x)(1944x^{3}-1569x^{2}+58x-1)\left(\frac{d}{dx}F(x)\right)^{2} \\ &+2F(x)(29376x^{3}+5580x^{2}-221x+1)\frac{d}{dx}F(x) \\ &+(144x^{2}-432x+1)F(x)^{2}=0 \end{aligned}$$